

$$\int_{n/4}^{m/6} \ln \Gamma(z) dz$$

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Abstract: We evaluate $\int \ln \Gamma(z) dz$, which could be notated $\psi_{-2}(z)$, between rational limits with denominators 1, 2, 3, 4, and 6. The simplest interesting result is (using $t!$ for $\Gamma(t+1)$)

$$\int_0^{1/2} \ln t! dt = \frac{\gamma + 3 \ln \pi}{8} - \frac{\ln 2}{6} - \frac{3\zeta'(2)}{4\pi^2} - \frac{1}{2} \approx -.04285374065.$$

We note that all integrals of the form $\int z^n \ln z! dz$ are expressible in terms of $\zeta'(-n, z)$, the first derivative of the Hurwitz zeta function.

Introduction

We establish

$$\begin{aligned} \int_0^{1/2} \ln t! dt &= \frac{\gamma + 3 \ln \pi}{8} - \frac{\ln 2}{6} - \frac{3\zeta'(2)}{4\pi^2} - \frac{1}{2} \approx -.04285374065, \\ \int_0^{1/3} \ln t! dt &= \frac{2\gamma + 5 \ln 2\pi}{18} - \frac{25 \ln 3}{72} - \frac{\sqrt{3}\pi}{54} + \frac{\sqrt{3}\psi_1(1/3)}{36\pi} - \frac{2\zeta'(2)}{3\pi^2} - \frac{1}{3} \approx -.02296612858214, \\ \int_0^{1/4} \ln t! dt &= \frac{3\gamma + 7 \ln \pi - 9 \ln 2}{32} - \frac{9\zeta'(2)}{16\pi^2} + \frac{\lambda}{4\pi} - \frac{1}{4} \approx -.014099944308, \\ \int_0^{1/6} \ln t! dt &= \frac{10\gamma - 23 \ln 3 + 22 \ln \pi}{144} - \frac{\sqrt{3}\pi}{36} + \frac{\sqrt{3}\psi_1(1/3)}{24\pi} - \frac{5\zeta'(2)}{12\pi^2} - \frac{1}{6} \approx -.006818642158464, \end{aligned}$$

where

$$\lambda := \text{Catalan's constant} = \frac{1}{8} \left(\psi_1\left(\frac{1}{4}\right) - \pi^2 \right) = \frac{1}{16} \left(\psi_1\left(\frac{1}{4}\right) - \psi_1\left(\frac{3}{4}\right) \right) \approx .9159655941772,$$

which, with the Γ multiplication formula and its limit, the Barnes-Alexeiewsky recurrence,

$$\int_{z-1}^z \ln t! dt = \ln \sqrt{2\pi} - (1 - \ln z)z,$$

suffice for $\int_{n/4}^{m/6}$. *E.g.*,

$$\begin{aligned} \int_0^{2/3} \ln t! dt &= \frac{\gamma + 10 \ln 2 + 4 \ln \pi}{9} - \frac{49 \ln 3}{72} + \frac{\sqrt{3}\pi}{54} - \frac{\sqrt{3}\psi_1(1/3)}{36\pi} - \frac{2\zeta'(2)}{3\pi^2} - \frac{2}{3} \approx -.06178124243647 \\ \int_0^{3/4} \ln t! dt &= \frac{3\gamma + 15 \ln \pi - 33 \ln 2}{32} + \frac{3 \ln 3}{4} - \frac{9\zeta'(2)}{16\pi^2} - \frac{\lambda}{4\pi} - \frac{3}{4} \approx -.06959909378, \\ \int_0^{5/6} \ln t! dt &= \frac{10\gamma - 119 \ln 3 + 70 \ln \pi}{144} + \frac{5 \ln 5 - 2 \ln 2}{6} + \frac{\sqrt{3}\pi}{36} - \frac{\sqrt{3}\psi_1(1/3)}{24\pi} - \frac{5\zeta'(2)}{12\pi^2} - \frac{5}{6} \\ &\approx -.07570064608869, \end{aligned}$$

The key lemma is an expansion of the general polylogarithm Li_k near 1.

Lemmas:

[Cohen, 1992] gives, without proof,

$$\text{Li}_k(e^{-z}) = \begin{cases} (-k)!z^{k-1} + \sum_{n \geq 0} \frac{\zeta(k-n)}{n!}(-z)^n, & \text{if } k \neq [k] \text{ or } k \leq 0; \text{ else} \\ (H_{k-1} - \ln z) \frac{(-z)^{k-1}}{(k-1)!} + \sum_{n \geq 0, n \neq k-1} \frac{\zeta(k-n)}{n!}(-z)^n, & H_{k-1} := \Psi_0(k) + \gamma; \end{cases} \quad (\text{Li1})$$

where

$$\text{Li}_k(z) := \sum_{n \geq 1} \frac{z^n}{n^k}.$$

([Gosper, 1990] gave (Li1) for negative integer k .) Differentiating ∂k at 1,

$$\begin{aligned} \sum_{n \geq 1} \frac{(\ln n)e^{-nz}}{n} &= - \left. \frac{\partial \text{Li}_k(e^{-z})}{\partial k} \right|_{k=1} \\ &= - \sum_{n \geq 1} \frac{\zeta'(1-n)}{n!}(-z)^n + \frac{(\gamma + \ln z)^2}{2} + \gamma_1 + \frac{\pi^2}{12}, \end{aligned} \quad (\text{Li1}')$$

where

$$\begin{aligned} \gamma_1 &= \frac{(\ln 2\pi)^2 - \gamma^2}{2} + \zeta''(0) + \frac{\pi^2}{24} \\ &= \sum_{n \geq 1} \frac{\ln n}{n} - \frac{(\ln(n+1))^2 - (\ln n)^2}{2} \approx -.07281584548367672, \end{aligned}$$

the first Stieltjes constant [Liang, 1972]. Differentiating the ζ reflection formula at a negative, even integer:

$$\zeta'(-2n) = (-)^n \frac{(2n)!\zeta(2n+1)}{2^{2n+1}\pi^{2n}}.$$

We can apply this to formula (Li1') if we first eliminate $\zeta'(1-2n)$ by setting $z = 2\pi it$ and taking the imaginary part. The resulting infinite series coincides with the Taylor expansion of $\ln \Gamma(t)$ at 0:

$$\begin{aligned} \ln \Gamma(t) &= \frac{1}{2} \left(\ln \Gamma(t) - \ln(-\Gamma(-t)) + \ln \frac{\pi}{t \sin \pi t} \right) \\ &= \frac{1}{2} \ln \frac{\pi}{t \sin \pi t} - \gamma t - \sum_{n \geq 1} \frac{\zeta(2n+1)}{2n+1} t^{2n+1} \\ &= \sum_{n \geq 1} \frac{\ln n}{\pi n} \sin 2\pi n t + \frac{1}{2} \left(\ln \frac{2\pi^2}{\sin \pi t} + \gamma \right) - (\gamma + \ln 2\pi)t, \quad 0 < t < 1, \end{aligned}$$

which is (the hard part of) the Fourier expansion of $\ln \Gamma(t)$. Finishing the expansion, via

$$\ln \sin \pi t = - \sum_{n \geq 1} \frac{\cos 2\pi n t}{n} - \ln 2, \quad 0 < t < 1, \quad (\ln \sin)$$

we get

$$\ln \Gamma(t) = \frac{\ln 2\pi}{2} + \sum_{n \geq 1} \frac{2(\gamma + \ln 2\pi n) \sin(2\pi n t) + \pi \cos(2\pi n t)}{2\pi n}, \quad 0 < t < 1. \quad (\text{Fou})$$

Differentiating (Li1) ∂k at 2 produces the same series as integrating equation (Li1') dz :

$$\begin{aligned} \sum_{n \geq 1} \frac{(\ln n) e^{-nz}}{n^2} &= - \left. \frac{\partial \text{Li}_k(e^{-z})}{\partial k} \right|_{k=2} \\ &= - \sum_{n \geq 2} \frac{\zeta'(2-n)}{n!} (-z)^n - \left(\frac{(\gamma + \ln z)(\gamma - 2 + \ln z)}{2} + \gamma_1 + \frac{\pi^2}{12} + 1 \right) z - \zeta'(2), \end{aligned}$$

Thus

$$\sum_{n \geq 1} \frac{\pi \sin(2\pi n t) - 2(\ln n) \cos(2\pi n t)}{4\pi^2 n^2} = \ln f_1(t-1) + (\gamma - 1 + \ln(2\pi)) \binom{t}{2} + \frac{\zeta'(2)}{2\pi^2}, \quad (\psi_{-2})$$

where

$$f_1(n) = 1^1 2^2 \dots n^n, \quad f_1(z) := (2\pi)^{-z/2} \exp \left(\binom{z+1}{2} + \int_0^z \ln t! dt \right),$$

which could equally well come from integrating (Fou) dt .

This series comes out in terms of $\zeta(2)$ and $\zeta'(2)$ for $t = j/4$ and $k/6$, yielding the identities asserted in the Introduction.

[Glaisher, 1895] denotes $\int \ln \Gamma(t) dt$ by ilg , but, if we can decide on the right constant of integration, it would seem logical to simply call it $\psi_{-2}(t)$, the first “negapolygamma”. By parts, it is equivalent to $\int t \psi_0(t) dt$. f_1 is sometimes called Barnes’s G function. $\zeta'(2)$ or simple functions thereof are sometimes called Glaisher’s or Porter’s constant [Knuth, 1976], [Pitteway, 1988].

It seems strange that, at rationals like $1/3$, $1/4$, and $1/6$, the function which interpolates $1^1 2^2 \dots n^n$, *i.e.*, $\int \ln \Gamma$ comes out in simpler constants, while the function which interpolates $1 \cdot 2 \cdot \dots n$, *i.e.*, Γ does not. However, if we denote $\ln \Gamma$ by ψ_{-1} , this extends for subscripts -1 and -2 a pattern where $\psi_n(j/4 \text{ or } k/6)$ simplifies iff n is even. What seems stranger is that the Fourier coefficients of $\ln \Gamma(t)$ should be elementary, while those of $\ln t$ are not.

Further results

(ψ_{-2}) can be repeatedly integrated by parts to give $\int_0^u t^n \ln \Gamma(t) dt$, but for $n > 0$, the only u for which the series has an obvious closed form are integers and half integers. *E.g.*,

$$\int_0^{1/2} t \ln t! = \frac{\gamma}{48} + \frac{\ln \pi}{12} - \frac{4\zeta'(2) + 7\zeta(3)}{32\pi^2} - \frac{1 + \ln 2}{16} \approx -.01317042784872,$$

$$\int_0^{1/2} t^2 \ln t! = \frac{\gamma + 5 \ln \pi}{192} - \frac{2\zeta'(2) + 3\zeta(3)}{32\pi^2} + \frac{15\zeta'(4)}{32\pi^4} - \frac{37 \ln 2}{1440} - \frac{1}{72} \approx -.00469460599089.$$

The earlier ψ_{-2} integrals, combined with the Γ and ψ reflection formulæ, yield $\int_{n/4}^{m/6} \ln \sin \pi t dt$. A superior alternative is to integrate equation $(\ln \sin)$,

$$\int_0^u \ln \sin \pi t dt = -u \ln 2 - \frac{1}{2\pi} \sum_{n \geq 1} \frac{\sin 2n\pi u}{n^2}, \quad 0 \leq u \leq 1,$$

which gives trigammmas for *all* rational u . *E.g.*,

$$\int_0^{1/5} \ln \sin \pi t dt = \frac{1}{25\pi} \left(\frac{4}{\sqrt{5}} \phi \pi^2 - \phi \psi_1 \left(\frac{1}{5} \right) - \psi_1 \left(\frac{2}{5} \right) \right) \sin \frac{\pi}{5} - \frac{\ln 2}{5} \approx -.2973633652639,$$

where

$$\phi := \frac{1 + \sqrt{5}}{2} \approx 1.6180339887499.$$

This, with the Γ reflection formula, gives $\ln \Gamma$ integrals over intervals symmetric about integers or half integers:

$$\int_{1-z}^z \ln t! dt = z \ln z - (1-z) \ln(1-z) + (1-2z)(1 - \ln \sqrt{\pi}) - \int_{1-z}^{1/2} \ln \sin \pi t dt.$$

E.g.,

$$\int_{2/5}^{3/5} \ln t! dt = \frac{\ln \pi - 3 \ln 2}{10} + \frac{1}{25\pi} \left(\phi \psi_1 \left(\frac{2}{5} \right) - \psi_1 \left(\frac{1}{5} \right) + \frac{2\pi^2}{\phi \sqrt{5}} \right) \sin \frac{\pi}{5} + \frac{3 \ln 3 - \ln 5 - 1}{5} \approx -.0238446115029.$$

With repeated integration of $(\ln \sin)$ by parts, we get $\int_0^u t^n \ln \sin \pi t dt$ in terms of polygammas for all positive integer n and rational u . *E.g.*,

$$\int_0^{1/8} t \ln \sin \pi t dt = \frac{\sqrt{2}}{2048\pi^2} \left(\psi_2 \left(\frac{1}{8} \right) - 4\psi_1 \left(\frac{1}{8} \right) \right) + \frac{5 + 4\sqrt{2}}{512} \pi + \frac{2\sqrt{2} - 1}{64\pi} \lambda + 7 \frac{37 + 16\sqrt{2}}{1024\pi^2} \zeta(3) - \frac{\ln 2}{128}$$

$$\approx -.01130943023.$$

The ψ_{-3} integral (equivalent to $\int t^2 \psi_0(t) dt$) over the general unit interval appears to require the ψ_{-2} integral over the general interval:

$$\int_{z-1}^z t \ln t! dt = \int_0^z \ln t! dt + \frac{z-2}{2} z \ln z - \frac{z^2}{4} + z - \frac{\ln 2\pi}{3} + \frac{\zeta'(2)}{2\pi^2}.$$

E.g., employing the $\psi_{-2}(1/6)$ result from the Introduction,

$$\int_{-5/6}^{1/6} t \ln t! dt = \frac{\sqrt{3}}{24\pi} \psi_1\left(\frac{1}{3}\right) - \frac{\gamma + 13 \ln 2\pi}{72} - \frac{1 + \ln 3}{144} + \frac{\zeta'(2)}{12\pi^2} \approx -.2815791335514.$$

The Hurwitz zeta function

All of these ψ_{-n} integrals are expressible via the first derivative of the Hurwitz zeta function $\zeta(s, a)$. Differentiating ∂s the reflection formula [Whittaker, 1927],

$$\frac{\partial}{\partial s} \frac{(2\pi)^s \zeta(1-s, a)}{\Gamma(s)} = - \sum_{n \geq 1} \frac{\pi \sin(s/2 - 2an)\pi + 2(\ln n) \cos(s/2 - 2an)\pi}{n^s},$$

which, for integer s , can be written

$$\zeta'(1-s, a) = B_s(a) \frac{\psi_0(s) - \ln 2\pi}{s} - (-)^s \frac{(s-1)!}{(2\pi)^{2s}} \frac{\partial^s}{\partial a^s} \sum_{n \geq 1} \frac{\pi \sin(2\pi an) - 2 \ln(n) \cos(2\pi an)}{n^{2s}},$$

where B_s is the Bernoulli polynomial

$$B_s(a) = -s\zeta(1-s, a).$$

With this we eliminate the series from equation (ψ_{-2}) ,

$$\int_0^a \ln \Gamma(t) dt = a \frac{1-a + \ln 2\pi}{2} + \zeta'(-1, a) - \zeta'(-1),$$

and from its integration by parts,

$$\int_0^a t \ln \Gamma(t) dt = \frac{a}{24} + a^2 \frac{1 + 2 \ln 2\pi}{8} - \frac{a^3}{4} + a \zeta'(-1, a) - \frac{\zeta'(-2, a)}{2} + \frac{\zeta(3)}{8\pi^2},$$

etc.

We also note that ζ' provides a concise, analytic extension of the “higher factorials”, *e.g.*:

$$1^{1^s} 2^{2^s} \dots n^{n^s} = e^{\zeta'(-s, n+1) - \zeta'(-s)},$$

$$1!2! \dots n! = n!^{n+1} e^{\zeta'(-1) - \zeta'(-1, n+1)} = e^{(n+1)(\zeta'(0, n+1) - \zeta'(0)) + \zeta'(-1) - \zeta'(-1, n+1)},$$

and reflection in the second argument provides polylogarithms, *e.g.*:

$$\text{Li}_2(e^{2i\pi z}) = \pi^2 B_2(z) + 2i\pi (\zeta'(-1, z) - \zeta'(-1, 1-z)).$$

With $z = 1/6$ and our earlier result for $\zeta'(-1, 1/6)$,

$$\text{Li}_2(e^{i\pi/3}) = \frac{\pi^2}{36} + \frac{i}{\sqrt{3}} \left(\frac{\psi_1(1/3)}{2} - \frac{\pi^2}{3} \right) \approx .2741556778080377 + 1.01494160640965i,$$

but this merely reminds us that we can express $\text{Li}_k(e^{ip\pi/q})$ in terms of $\psi_{k-1}(j/q)$, $0 < j < q$, simply by grouping the terms of the Li series q at a time, leading to a rational function summand.

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